

---

*Transverse polarization of  $\Lambda^0$  hyperons  
in quasi-real photoproduction:  
Quark Recombination Model*

I. ALIKHANOV

Saint Petersburg State University, Russia

O. GREBENYUK

Petersburg Nuclear Physics Institute, Russia

---

*XII Workshop on High Energy Spin Physics*

*Dubna, September 3-7, 2007*

---

---

## OUTLINE

- ✓ Introduction
- ✓ Experimental Field
- ✓ Quark Scattering Model (QSM)
- ✓ Quark Recombination Model (QRM)
- ✓ Calculations and Results
- ✓ Conclusion

---

## INTRODUCTION

- ✓ Among the hyperons,  $\Lambda^0$  fills a special place due to the spin-flavor structure of its wave function within the SU(6) symmetry

$$|\Lambda\rangle_{\frac{1}{2}} = |ud\rangle_0 |s\rangle_{\frac{1}{2}}.$$

- ✓ In unpolarized reactions  $ab \rightarrow \Lambda X$ , the direction of the polarization is defined by

$$\mathbf{n} \propto [\mathbf{p}_a \times \mathbf{p}_\Lambda].$$

- ✓ The polarization sign

$$\mathbf{P} \cdot \mathbf{n} < 0 \quad pp \rightarrow \Lambda X,$$

$$\mathbf{P} \cdot \mathbf{n} > 0 \quad K^- p \rightarrow \Lambda X.$$

---

## EXPERIMENTAL FIELD

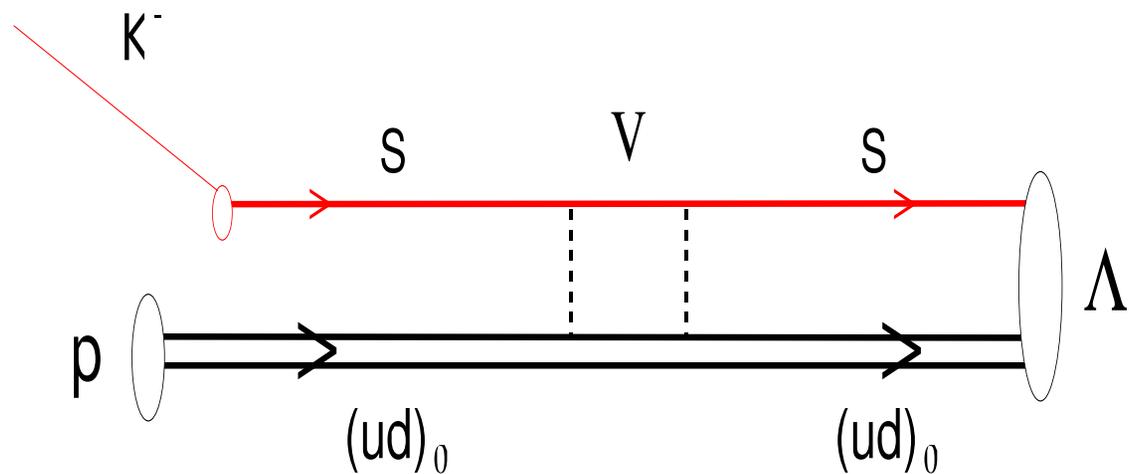
The recent HERMES results [[arXiv:0704.3133](https://arxiv.org/abs/0704.3133)] are qualitatively similar to the polarization in  $K^-p$ .

- ✓ the positive sign has been observed in both the reactions.
- ✓ similarity in the  $p_T$  dependence.

In the current fragmentation region ( $x_F > 0$ ), the  $\Lambda$  kinematic is mostly determined by the strange quark.

---

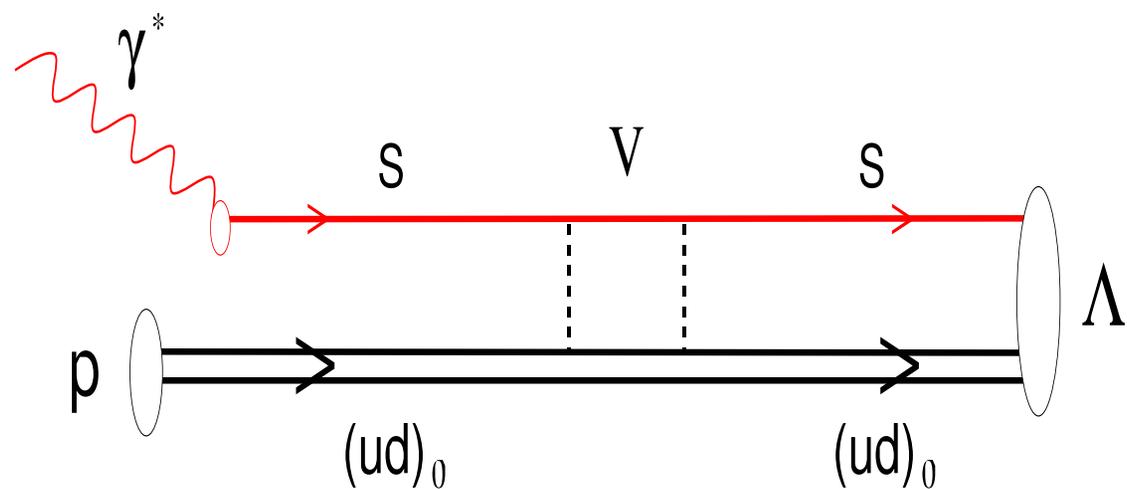
## EXPERIMENTAL FIELD



We consider the  $x_F > 0$  region only.

---

## EXPERIMENTAL FIELD



We consider the  $x_F > 0$  region only.



---

## QSM FOR THE $\Lambda$ PHOTOPRODUCTION

Since the polarization at HERMES is available versus  $\zeta$  and  $p_T$ , the QSM has been rewritten in terms of the variables

$$\zeta_{i(f)} = \frac{E_{i(f)} + p_{zi(f)}}{E_b + p_{zb}},$$

$$P\left(\frac{\zeta_f}{\zeta_i}, p_T\right) = -\frac{2C\alpha_s V}{1 + V^2 \cos^2 \theta/2} \frac{\sin^3 \theta/2 \ln(\sin \theta/2)}{\cos \theta/2},$$

$$V = V\left(\frac{\zeta_f}{\zeta_i}, p_T\right), \quad \theta = \theta\left(\frac{\zeta_f}{\zeta_i}, p_T\right)$$

---

## QSM FOR THE $\Lambda$ PHOTOPRODUCTION

We consider the  $\zeta > 0.25$  region only, which presumably relates to the current fragmentation.

One needs to know the  $\zeta_i$  as well as the  $\zeta_f$  distributions.

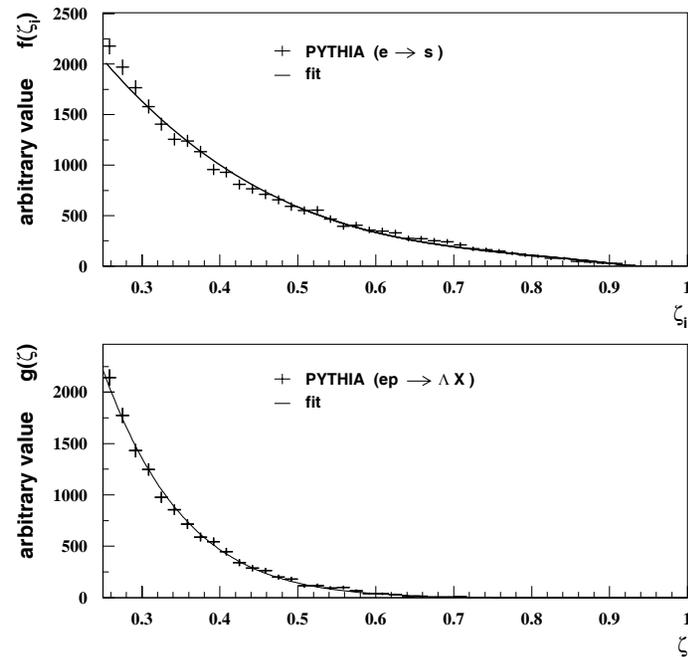
$$\zeta_f = \frac{m_s}{m_\Lambda} \zeta, \quad p_T = \frac{m_s}{m_\Lambda} p_T.$$

$$P_\zeta = \int d\zeta_i dp_T h(p_T) P\left(\frac{\zeta}{\zeta_i}, p_T\right) f(\zeta_i),$$

$$P_{p_T} = \int d\zeta_i d\zeta g(\zeta) P\left(\frac{\zeta}{\zeta_i}, p_T\right) f(\zeta_i).$$

---

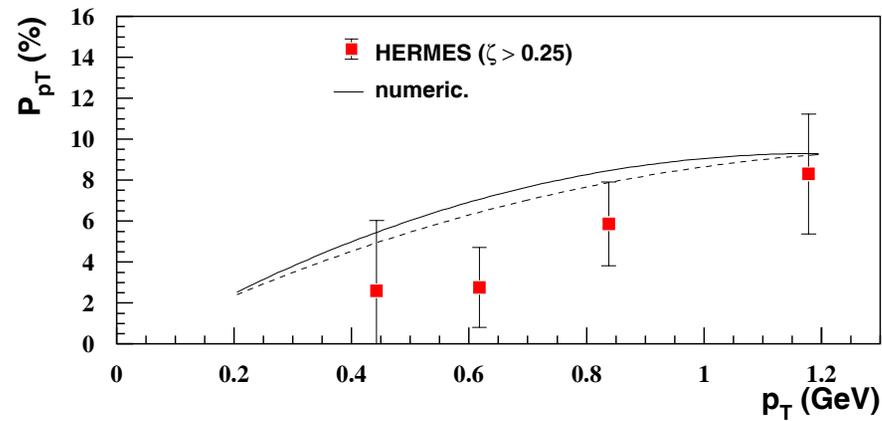
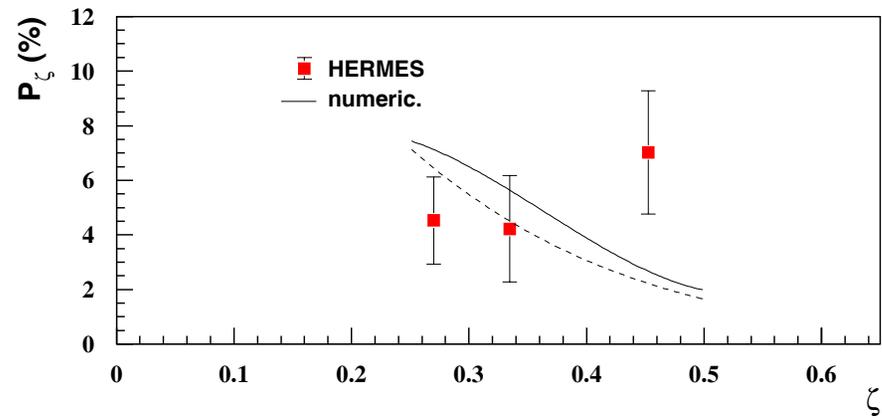
# CALCULATIONS



$$h(p_T) \propto \exp(-4.2p_T^2), \quad [\text{Acta. Phys. Polon., B33(2002)3785}].$$



# RESULTS



---

## QUARK RECOMBINATION MODEL

Y. Yamamoto, K. Kubo and H. Toki [[Prog. Theor. Phys. 98\(1997\)95](#)].

The polarization is standardly given by

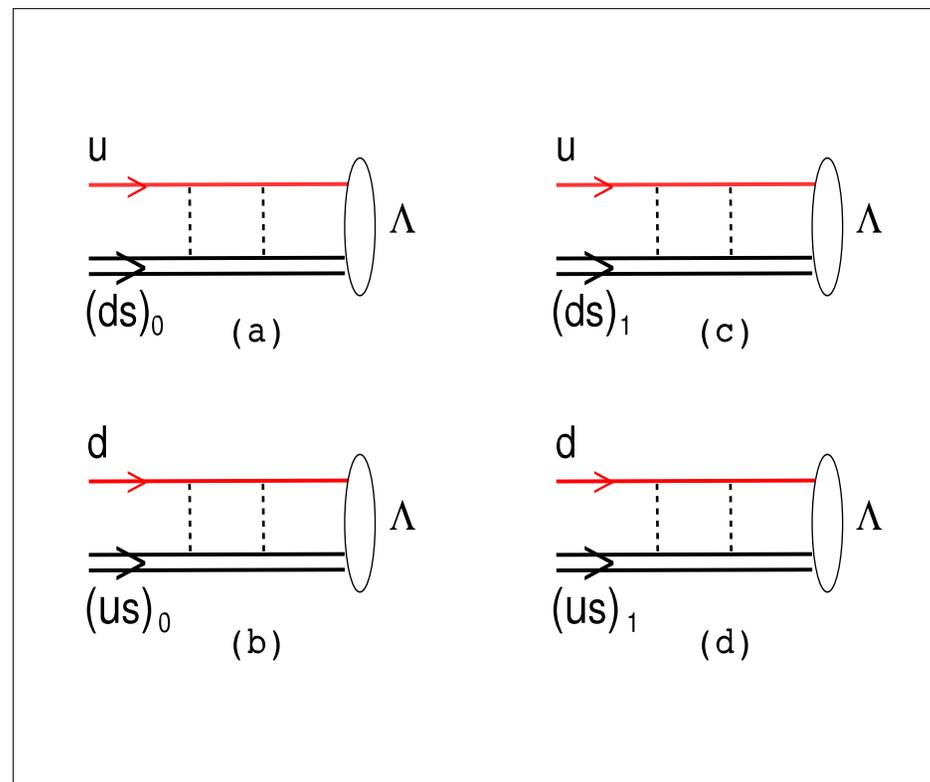
$$P = \frac{\sum_{M_i} |\langle +1/2 | S | M_i \rangle|^2 - \sum_{M_i} |\langle -1/2 | S | M_i \rangle|^2}{\sum_{M_i} |\langle +1/2 | S | M_i \rangle|^2 + \sum_{M_i} |\langle -1/2 | S | M_i \rangle|^2}.$$

$$|\langle M_f | S | M_i \rangle|^2 = \int G^{M_f}(r_f) |M(r_f, r_i)|^2 G^{M_i}(r_i).$$

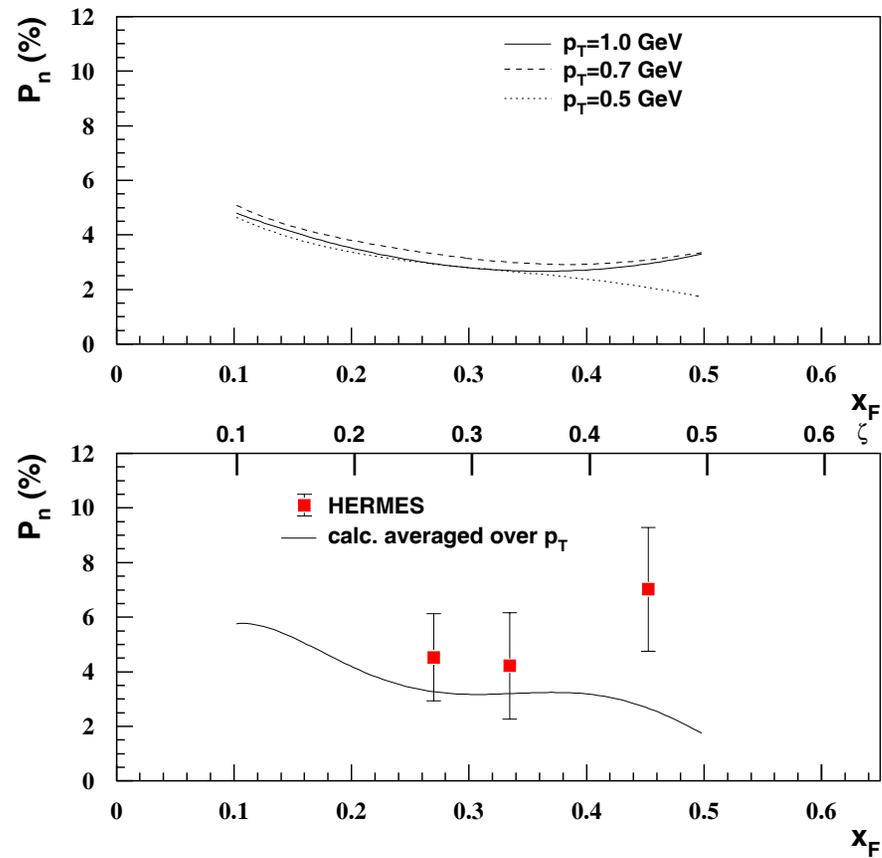
The interaction is assumed to be scalar.

## QRM FOR THE $\Lambda$ PHOTOPRODUCTION

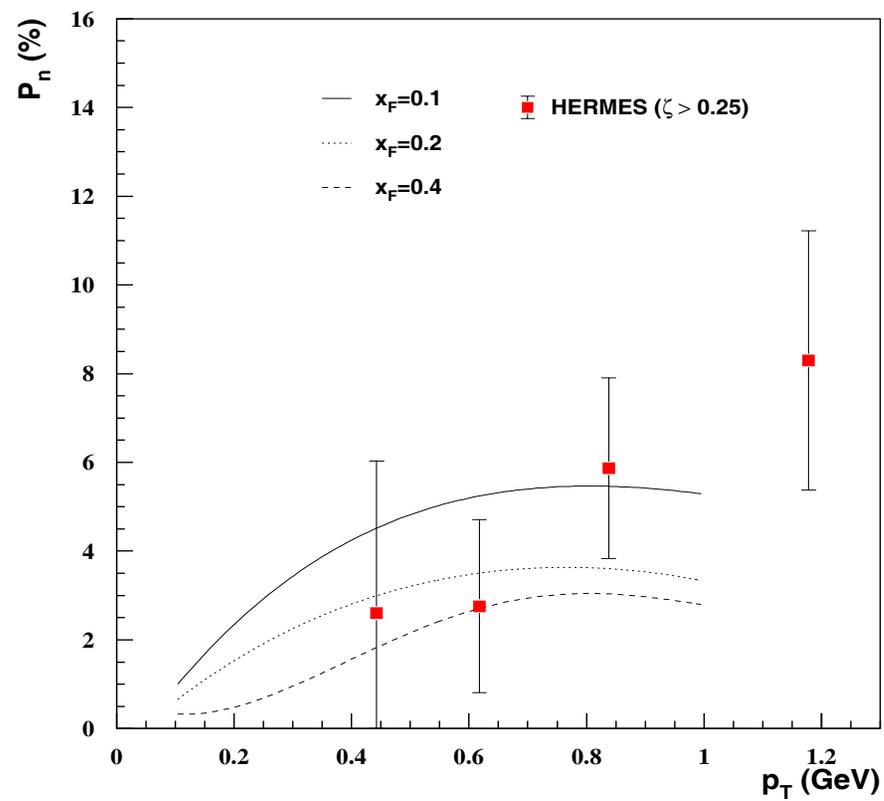
The photoproduction may be fairly expected to be richer with the subprocesses



# RESULTS



# RESULTS



---

## CONCLUSION

- ✓ The reached reproduction should be regarded only as qualitative
- ✓ The calculations are based on the SU(6) symmetry, while it is not exact
- ✓ We used the PYTHIA programm, which gives rather qualitative than quantitative predictions
- ✓ We assumed the  $s+(ud)_0$  to take place only
- ✓ A large difficulty is the parameter  $2C\alpha_s$

---

Thank You!